

ME 201/APh 250, Homework 2:

Assigned: Friday, Apr 12, 2019

Due: Friday, Apr 19, 2019

N&C 4.8:

Let us expand the unitary $U = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ in terms of the set of Pauli matrices $\{I, X, Y, Z\}$ (note this set forms a complete basis set of all 2x2 Hermitian matrices). We let $U = a_0I + a_1X + a_2Y + a_3Z$ where

$$a_0 = \frac{a+d}{2}, a_1 = \frac{b+c}{2}, a_2 = \frac{c-b}{2i}, a_3 = \frac{a-d}{2}$$

Since U is unitary, $U^\dagger U = I$. Then, we have the following list of constraints for a_0, a_1, a_2, a_3

$$\begin{aligned} \|a_0\|^2 + \|a_1\|^2 + \|a_2\|^2 + \|a_3\|^2 &= 1 \\ a_0^* a_1 + a_1^* a_0 + i a_2^* a_3 - i a_3^* a_2 &= 0 \\ a_0^* a_2 - i a_1^* a_3 + a_2^* a_0 + i a_3^* a_1 &= 0 \\ a_0^* a_3 + i a_1^* a_2 - i a_2^* a_1 + a_3^* a_0 &= 0 \end{aligned}$$

Let $\cos(\frac{\theta}{2}) = \|a_0\|$, from the first constraint, we get $\|a_1\|^2 + \|a_2\|^2 + \|a_3\|^2 = \|\sin(\frac{\theta}{2})\|$ Let

$$\begin{aligned} n_x &= \|a_1\|/\|\sin(\theta/2)\| \\ n_y &= \|a_2\|/\|\sin(\theta/2)\| \\ n_z &= \|a_3\|/\|\sin(\theta/2)\| \end{aligned}$$

Our definition satisfies $n_x^2 + n_y^2 + n_z^2 = 1$. We now define the following:

$$\begin{aligned} a_0 &= e^{i\alpha} \cos(\theta/2) \\ a_1 &= -ie^{i\alpha} \sin(\theta/2) n_x \\ a_2 &= -ie^{i\alpha} \sin(\theta/2) n_y \\ a_3 &= -ie^{i\alpha} \sin(\theta/2) n_z \end{aligned}$$

It can be shown these satisfies the remaining constraints. Therefore,

$$\begin{aligned} U &= a_0I + a_1X + a_2Y + a_3Z \\ &= e^{i\alpha} (\cos(\theta/2)I - i\sin(\theta/2)(n_xX + n_yY + n_zZ)) \\ &= e^{i\alpha} R_{\vec{n}}(\theta) \end{aligned}$$

N&C 4.17:

We can check that $CNOT = (I \otimes H)(CZ)(I \otimes H)$

N&C 4.32:

In the discussion in **N&C 2.4.1** from page 99 to 101, throwing away information after measurement results in a density matrix given by:

$$\rho' = \sum_m M_m \rho M_m^\dagger$$

Since we are performing projective measurements, it turns out that $M_m = P_m, M_m^\dagger = P_m^\dagger = P_m = M_m$. Then,

$$\rho' = P_0 \rho P_0 + P_1 \rho P_1$$

To show that $\text{tr}_2(\rho) = \text{tr}_2(\rho')$, consider $\rho = \sum_i p_i |\psi_i\rangle \langle \psi_i|, |\psi_i\rangle = \sum_{mn} c_i^{mn} |mn\rangle$,

$$\begin{aligned} \text{tr}_2 \rho &= \text{tr}_2 \left(\sum_i p_i \sum_{mnr s} c_i^{mn} c_i^{rs*} |mn\rangle \langle rs| \right) \\ &= \sum_i p_i \sum_{mnr s} c_i^{mn} c_i^{rs*} |m\rangle \langle r| \text{tr}(|n\rangle \langle s|) \end{aligned}$$

$$\begin{aligned} \text{tr} \rho' &= \text{tr}_2 \left(\sum_i p_i \sum_{mnr s} c_i^{mn} c_i^{rs*} |m\rangle \langle r| (\langle 0|n\rangle \langle s|0\rangle |0\rangle \langle 0| + \langle 0|n\rangle \langle s|1\rangle |0\rangle \langle 1| \right. \\ &\quad \left. + \langle 1|n\rangle \langle s|0\rangle |1\rangle \langle 0| + \langle 1|n\rangle \langle s|1\rangle |1\rangle \langle 1|) \right) \\ &= \text{tr}_2 \left(\sum_i p_i \sum_{mnr s} c_i^{mn} c_i^{rs*} |m\rangle \langle r| [\langle 0|n\rangle \langle s|0\rangle + \langle 1|n\rangle \langle s|1\rangle] \right) \\ &= \sum_i p_i \sum_{mnr s} c_i^{mn} c_i^{rs*} |m\rangle \langle r| \text{tr}(|n\rangle \langle s|) \end{aligned}$$

N&C 4.33:

The circuit is given by

$$\begin{aligned} C &= \frac{1}{\sqrt{2}} \left(\begin{bmatrix} 1 & 1 \\ 1 & -1 \end{bmatrix} \otimes \begin{bmatrix} 1 & 0 \\ 0 & 1 \end{bmatrix} \right) \begin{bmatrix} 1 & 0 & 0 & 0 \\ 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & 1 \\ 0 & 0 & 1 & 0 \end{bmatrix} \\ &= \frac{1}{\sqrt{2}} \begin{bmatrix} 1 & 0 & 0 & 1 \\ 0 & 1 & 1 & 0 \\ 1 & 0 & 0 & -1 \\ 0 & 1 & -1 & 0 \end{bmatrix} \end{aligned}$$

The action of the circuit on the bell state B_{00} is

$$\begin{aligned} CB_{00} &= C \frac{1}{\sqrt{2}} \begin{bmatrix} 1 \\ 0 \\ 0 \\ 1 \end{bmatrix} \\ &= \begin{bmatrix} 1 \\ 0 \\ 0 \\ 0 \end{bmatrix} \\ &= |00\rangle \end{aligned}$$

Similarly,

$$CB_{01} = |01\rangle, CB_{10} = |10\rangle, CB_{11} = |11\rangle$$

It is clear our measurement operators are given as:

$$\begin{aligned} M_1 &= |00\rangle \langle B_{00}| \\ M_2 &= |01\rangle \langle B_{01}| \\ M_3 &= |10\rangle \langle B_{10}| \\ M_4 &= |11\rangle \langle B_{11}| \end{aligned}$$

Then, the POVM elements are:

$$\begin{aligned} E_1 &= M_1^\dagger M_1 = |B_{00}\rangle \langle B_{00}| \\ E_2 &= M_2^\dagger M_2 = |B_{01}\rangle \langle B_{01}| \\ E_3 &= M_3^\dagger M_3 = |B_{10}\rangle \langle B_{10}| \\ E_4 &= M_4^\dagger M_4 = |B_{11}\rangle \langle B_{11}| \end{aligned}$$

They are the projectors form from your Bell states.

Qns 2:

We note that the total system AB is spanned by the orthonormal basis vectors: $\{|00\rangle, |01\rangle, |10\rangle, |11\rangle\}$. In this basis, the wave function can be rewritten as $|\psi\rangle = u_1 |00\rangle + u_2 |01\rangle + u_3 |10\rangle + u_4 |11\rangle$. Then, the density matrix is given by:

$$\begin{aligned} \rho &= |\psi\rangle \langle \psi| \\ &= (u_1 |00\rangle + u_2 |01\rangle + u_3 |10\rangle + u_4 |11\rangle)(u_1^* \langle 00| + u_2^* \langle 01| + u_3^* \langle 10| + u_4^* \langle 11|) \\ &= u_1 u_1^* |00\rangle \langle 00| + u_1 u_2^* |00\rangle \langle 01| + u_1 u_3^* |00\rangle \langle 10| + u_1 u_4^* |00\rangle \langle 11| \\ &\quad + u_2 u_1^* |01\rangle \langle 00| + u_2 u_2^* |01\rangle \langle 01| + u_2 u_3^* |01\rangle \langle 10| + u_2 u_4^* |01\rangle \langle 11| \\ &\quad + u_3 u_1^* |10\rangle \langle 00| + u_3 u_2^* |10\rangle \langle 01| + u_3 u_3^* |10\rangle \langle 10| + u_3 u_4^* |10\rangle \langle 11| \\ &\quad + u_4 u_1^* |11\rangle \langle 00| + u_4 u_2^* |11\rangle \langle 01| + u_4 u_3^* |11\rangle \langle 10| + u_4 u_4^* |11\rangle \langle 11| \end{aligned}$$

We note that $\text{Tr}_B(|pq\rangle \langle rs|) = \langle q|s\rangle |p\rangle \langle r|$. We note that the partial trace is an inner operator. Therefore, taking partial trace over B, we obtain:

$$\begin{aligned} \text{Tr}_B(\rho) &= u_1 u_1^* |0\rangle \langle 0| + u_1 u_3^* |0\rangle \langle 1| \\ &\quad + u_2 u_2^* |0\rangle \langle 0| + u_2 u_4^* |0\rangle \langle 1| \\ &\quad + u_3 u_1^* |1\rangle \langle 0| + u_3 u_3^* |1\rangle \langle 1| \\ &\quad + u_4 u_2^* |1\rangle \langle 0| + u_4 u_4^* |1\rangle \langle 1| \end{aligned}$$

Qns 3:

Note that $H|0\rangle = \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle)$ and $H|1\rangle = \frac{1}{\sqrt{2}}(|0\rangle - |1\rangle)$. Therefore, $H|x\rangle = \frac{1}{\sqrt{2}}(|0\rangle + (-1)^x|1\rangle)$ where $x \in \{0, 1\}$. We can rewrite this as $H|x\rangle = \frac{1}{\sqrt{2}} \sum_{z \in \{0,1\}} (-1)^{xz} |z\rangle$. Then,

$$\begin{aligned}
H^{\otimes n} |x_1\rangle |x_2\rangle \dots |x_n\rangle &= (H|x_1\rangle)(H|x_2\rangle)\dots(H|x_n\rangle) \\
&= \frac{1}{\sqrt{2}} \sum_{z_1 \in \{0,1\}} (-1)^{x_1 z_1} |z_1\rangle \frac{1}{\sqrt{2}} \sum_{z_2 \in \{0,1\}} (-1)^{x_2 z_2} |z_2\rangle \dots \frac{1}{\sqrt{2}} \sum_{z_n \in \{0,1\}} (-1)^{x_n z_n} |z_n\rangle \\
&= \frac{1}{\sqrt{2^n}} \sum_{z_1, z_2, \dots, z_n \in \{0,1\}} (-1)^{x_1 z_1 + x_2 z_2 + \dots + x_n z_n} |z_1\rangle |z_2\rangle \dots |z_n\rangle
\end{aligned}$$